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## Predicting the Impact Sensitivities of Energetic Materials through Zone-Center Phonon Up-Pumping

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**Abstract:** The development of new energetic materials (EMs) is accompanied by significant hazards, prompting interest in their computational design. Before reliable *in silico* design strategies can be realized, however, approaches to understand and predict EM response to mechanical impact must be developed. We present here a fully *ab initio* model based in phonon up-pumping which successfully ranks the relative impact sensitivity of a series of organic EMs. The methodology depends only on the crystallographic unit cell and Brillouin zone center vibrational frequencies. We therefore expect this approach to become an integral tool in the large-scale screening of potential EMs.

The mechanically induced reactivity of solids represents a fascinating and poorly understood aspect of materials science. With respect to energetic materials (explosives, propellants, pyrotechnics, and gas generators; EMs) this reactivity is crucial for determining their handling safety and potential for real-world applications in both civilian and military technologies.<sup>1</sup> EMs are characterized by their ability to rapidly convert large amounts of chemical potential energy into kinetic energy, with potentially devastating effects if uncontrolled. Owing to stringent safety and performance requirements (e.g. detonation pressure and velocity, and the heat of explosion), only a limited number of EMs have found widespread application. However, many of the well-established EMs do not comply with increasing environmental and public health regulations, prompting the need to develop new, non-toxic and ‘green’ EMs which neither compromise on safety nor on performance.<sup>2</sup>

The impact sensitivity (IS) of an EM represents an internationally recognized material hazard classification criterion.<sup>3</sup> Where materials are highly sensitive (*i.e.* initiate on mild impact), they are generally not fit for real-world applications, regardless of their other performance characteristics. Correspondingly, significant efforts have been devoted to identifying strategies for characterizing, rationalizing, and designing low sensitivity EMs (LSEMs) which offer greater potential for safe handling. Various experimental protocols exist to measure the IS of EMs.<sup>4</sup> These approaches typically involve the use of a drop hammer device, such as the BAM Fall-Hammer or the Rotter Impact device, both requiring the preparation of bulk quantities of novel EMs. Thus, while the discovery of new EMs requires significant investment of both time and financial resources, the lack of *a priori* knowledge regarding their explosive properties also carries considerable risk to health and safety.

The challenges presented by the experimental discovery of LSEMs have prompted efforts to derive theoretical descriptions of EM sensitivity.<sup>5,6</sup> Many efforts have focused on the screening of molecular parameters. For example, bond dissociation energies have been successfully correlated against EM sensitivity,<sup>7</sup> including through machine learning analyses.<sup>8,9</sup> A correlation between IS and molecular oxidizing potential (*i.e.* oxygen balance) has been successfully demonstrated for a series of polynitroaromatic explosives,<sup>10</sup> although attempts to verify the same correlation for a series of PETN analogues failed.<sup>11</sup> Other molecule-centric approaches have included the correlation of sensitivity with

NMR chemical shifts,<sup>12</sup> atomic composition and their corresponding HOMO-LUMO gaps,<sup>13</sup> electronegativity,<sup>14</sup> and charge density distribution.<sup>15</sup>

While these approaches have proved effective for rationalizing sensitivity trends in chemically related series of EMs, many important phenomena remain outside their predictive capacity. Reports have shown, for instance, that IS can vary for the same EM expressed in different polymorphic forms. For example, HMX, (octahydro-1,3,5,7-tetranitro-1,3,5,7-tetrazocine, see Figure 1) exists in the thermodynamically stable  $\beta$ -form under ambient conditions, but it is often contaminated with traces of the  $\alpha$ - or  $\gamma$ -forms which are both more sensitive to mechanical impact.<sup>16,17</sup> An even more sensitive form ( $\delta$ -HMX) is readily accessible upon heating.<sup>18</sup> Similarly, a growing number of studies have shown that the formation of multi-component crystals (salts and co-crystals) can have significant influence on IS properties.<sup>19,20</sup> The crystal lattice therefore must play a role in determining IS, and consequently crystal-centric approaches have become the focus of growing attention. To this end IS correlations have been sought against a number of solid-state characteristics, including intermolecular bond distances,<sup>21</sup> electronic band gaps,<sup>22</sup> crystal void space,<sup>23,24</sup> and crystal morphology.<sup>25</sup> However, these approaches generally lack physical insight into the initiation mechanism and have therefore been largely engulfed by more sophisticated models. For example, identifying the need to associate bulk compressibility with electronic excitation, Bondarchuk successfully correlated IS against pressure-induced metallization for metal azides<sup>26</sup> and across a breadth of molecular EMs.<sup>27</sup> In an alternative, albeit related, approach various authors considered the development of local stress in perturbed solids, which results from anisotropic distortion of the solid under mechanical loading.<sup>28,29</sup> Following on from these efforts, correlations between crystal packing geometry (e.g.  $\pi \cdots \pi$  stacking)<sup>30,31</sup> have become increasingly popular.<sup>28,32</sup>

Following the pioneering work of Dlott and co-workers,<sup>33</sup> we recently presented an *ab initio* phonon up-pumping based approach for predicting EM IS.<sup>34,35</sup> This up-pumping model captures critical features of the impact event, including compression-induced excitation<sup>36</sup> and vibrationally induced metallization.<sup>37</sup> Our previous model was based on computationally expensive phonon dispersion curves (PDCs). The significant time and resources required to obtain PDCs represents a substantial barrier to implementing our model for larger-scale material screening. Moreover, the need for PDCs places large energetic molecules in low symmetry unit cells out of reach. A recent study by Bernstein<sup>38</sup> suggested that up-pumping based approaches using limited sampling of the Brillouin zone can be sufficient to achieve positive correlation. This has motivated our current study to explore the potential to expedite our computational screening approach, by considering only the Brillouin zone-center ( $\Gamma$ -point) *ab initio* vibrational frequencies. A key consideration in our up-pumping model is the assignment of  $\Omega_{max}$ , the vibrational frequency denoting the change in mode character from lattice to molecular behavior. In this work we explore assignment through three separate methods, namely mode counting, consideration of the vibrational mode contribution to the heat capacity and finally through tracking the changes in the molecular center of mass expressed by the eigenvectors. Herein we briefly outline the theoretical background of our phonon up-pumping model, and apply it to a well-known structurally and energetically diverse set of EMs.

## Theory

The mechanical compression of a material leads to a bulk increase in the thermodynamic energy  $\Delta U$ . This energy distributes between heat  $H$  and work  $W$ , such that  $\Delta U = H + W$ . Partition of the input mechanical energy, and hence the elevation of heat in the system, can be approximated from consideration of the compression factor  $V/V_0$  and the Grüneisen parameter,  $\gamma$ . The final temperature  $T_f$  that is reached by an

adiabatically compressed solid at initial temperature  $T_0$  with initial volume  $V_0$  can be described according to<sup>40</sup>

$$\left(\frac{T_f}{T_0}\right) = \left(\frac{V}{V_0}\right)^{-\gamma} \quad \text{Eqn 1}$$

The total heat transferred to the system can be re-written in terms of the bulk material heat capacity,  $C(T)$ , with

$$H = \int_{T_0}^{T_f} dT C(T) \quad \text{Eqn 2}$$

Upon mechanical impact, a compressive wave passes through the material, coupling most strongly to the lattice acoustic modes. Impact energy therefore transfers into these few branches.<sup>41</sup> Owing to the significant phonon-phonon scattering that occurs between lattice phonon modes, this energy rapidly redistributes throughout the external acoustic and optical branches. A good approximation for the initial conditions directly following mechanical impact is therefore to assume that the energy from the impact event is distributed across all of the lattice modes, together referred to as the phonon bath.<sup>33</sup> The initial phonon bath quasi-temperature  $\phi_{ph}$  condition can then be described by replacing  $C(T)$  with the phonon heat capacity  $C_{ph}$ , such that

$$C_{ph}[\phi_{ph} - T_0] = \int_{T_0}^{T_f} dT C(T) \quad \text{Eqn 3}$$

assuming that the phonon quasi-temperature must be a state of the bulk temperature-volume curve.<sup>33</sup> As the phonon bath represents only a small number of the total vibrational states in the material, the phonon heat capacity is considerably lower than the bulk heat capacity. Correspondingly, the phonon quasi-temperatures reach significantly higher values than will be achieved after thermalization of the material. Hence, the initial conditions for our model comprise a vibrationally 'hot' phonon region, and a vibrationally 'cold' molecular region of vibrational states.

An EM initiation event results from the rupture of a covalent bond. The energy stored in the phonon bands must therefore up-convert to reach an active molecular vibrational state, which is termed a target mode. To this end we consider the vibrational Hamiltonian of a molecular crystal,

$$H = H_i + H_L + H_c \quad \text{Eqn 4}$$

where  $H_i$  describes the internal motion of molecules,  $H_L$  denotes the motion of molecules about their equilibrium positions (*i.e.* the lattice vibrations), and  $H_c$  provides a route for coupling between the internal and external modes. For the general case of indirect phonon up-pumping,  $H_c$  can be expanded as,

$$H_c \propto \sum_{i,j,k} |V_{i,j,k}^3| \delta(\Omega_i - \omega_j - \omega_k) \quad \text{Eqn 5}$$

where

$$V_{i,j,k}^3 = \frac{\partial^3 E}{\partial u_i \partial u_j \partial u_k}$$

For simplicity, we have assumed Einstein phonons in Equation 5, thereby avoiding explicit consideration of wavevector dispersion. This is generally a good approximation for the internal vibrations of molecular

materials. The Kroneker delta of Equation 5, which ensures momentum and energy conservation, is conveniently rewritten as  $\rho^{(2)}$ . This is the two-phonon density of states and describes the formation of a new phonon state  $\Omega_i$  from the scattering of two lower energy phonons  $\omega_j$  and  $\omega_k$ . The coefficient  $V^3$  describes the anharmonic coupling strength between phonons  $i, j$  and  $k$ . Spectroscopic evidence<sup>42</sup> and semi-empirical calculations<sup>43</sup> have suggested that the values of  $V^3$  are approximately constant for molecular energetic materials. Correspondingly, the relative scattering depends only on the function  $\rho^{(2)}$ . Early developments<sup>33</sup> demonstrated that vibrational up-pumping follows a multi-staged approach. Energy from the ‘hot’ phonon bath (*i.e.* from  $\omega < \Omega_{max}$ , where  $\Omega_{max}$  denotes the top of the phonon bath) is first transferred to the doorway states (denoted by  $\Omega_{max} < \omega < 2\Omega_{max}$ ) with a rate given by  $\tau_1$ .

$$\tau_1 = \rho^{(2)}(\omega = \omega_{T/2}) \times [n_{ph} - n_d] \quad \text{Eqn 6a}$$

$$\tau_2 = \rho^{(2)} \times [n_{ph} + n_d - n_T] \quad \text{Eqn 6b}$$

Here it is assumed that  $\tau_1$  transfers all excess population energy (denoted by the Bose-Einstein populations,  $n$ ) from the phonon states ( $n_{ph}$ ) into the doorway region ( $n_d$ ), thereby re-establishing equilibrium populations of  $n_{ph}$  prior to further up-pumping. The ‘hot’  $n_d$  states subsequently up-pump to the target modes ( $n_T$ ) in the second step according to rate  $\tau_2$ . Where dispersion of  $\omega(\mathbf{k})$  is large (*i.e.* in the phonon bath), the Einstein approximation ( $\omega(\mathbf{k}) = \omega$ ) is not suitable. This leads to the energy restriction in Equation 6a. Where the target mode ( $n_T$ ) is adequately described as an Einstein phonon, explicit consideration of wavevector conservation can be neglected. This is done in Equation 6b. For simplicity we can borrow the terms *overtone* and *combination* from vibrational spectroscopy to describe the analogous processes involved in Equations 6a and 6b, respectively. In this terminology, we note that the contributions to  $\rho^{(2)}$  in Equation 6a are limited to the first overtone, as described by  $\omega = \omega_{T/2}$ .

All up-pumping steps are assumed to require delocalized vibrations - *i.e.* to include at least one phonon bath mode. Correspondingly, we generate  $\rho^{(2)}$  in Equations 6 by assuming at least one phonon must fall at  $\omega < \Omega_{max}$ . Similarly, the second phonon mode must not be in thermal equilibrium with the bulk and must hence fall at  $\omega < 2\Omega_{max}$  in Equation 6b. The highest available excitation by our phonon up-pumping model is therefore  $3\Omega_{max}$ .

The up-pumping model developed here is reminiscent of the starvation kinetics model developed by Eyring for unimolecular reactions at a shock front.<sup>39</sup> In the Eyring model, a dissociating bond (*i.e.* the target bond) is at equilibrium with a ‘cold’ (or ‘starved’) reservoir of oscillators (analogous to doorway modes in our model). As the shock front propagates through the material, it does not equilibrate directly with the dissociating bond, but rather transfers energy to the starved reservoir, which in turn ‘feeds’ the dissociating bond (*c.f.* our two-step up-pumping mechanism). Although the conceptual construction of our model is consistent with the Eyring model, the formalisms differ. Most notably, the Eyring model derives its system dependence by considering bond dissociation energies, thereby requiring knowledge of the primary decomposition event. In contrast, system dependence in our model follows from the relative vibrational energy transfer pathways, without any need for *a priori* knowledge of the material reactivity. We envision, however, future need to include into our model concepts developed by Eyring.

## Computational Details

*Gas phase simulations.* All gas phase simulations were performed in ORCA v4.2.<sup>44</sup> The molecular structures were extracted from available crystallographic data, and fully relaxed at the PBE/def2-TZVP level. The force constant matrix was calculated analytically for each relaxed geometry. Owing to the known deficiency of this level of theory at reproducing experimental frequencies, we adopt the corresponding scaling constant 1.0306.<sup>45</sup>

*Simulations of periodic systems.* All periodic DFT calculations were performed within the CASTEP suite v 19.11.<sup>46</sup> All input structures were taken from the Cambridge Crystallography Data Centre (CCDC) database: Hexanitrobenzene (REF: HNOBEN),  $\epsilon$ -CL-20 (REF: PUBMUU12),  $\beta$ -HMX (REF: OCHTET13), m-TNT (REF: ZZMUC),  $\alpha$ -FOX-7 (REF: SEDTUQ01),  $\alpha$ -Nitrotriazolone (REF: QOYJOD06), TATB (REF: TATNBZ03). In each case, the electronic structure was expanded in plane waves to a maximum cut-off energy of 1200 eV, and the nuclear-electron interactions were approximated with norm-conserving pseudopotentials as generated on-the-fly in CASTEP. The Brillouin zone was sampled on a Monkhorst-Pack grid<sup>47</sup> with spacings no greater than  $0.05 \text{ \AA}^{-1}$ . Owing to its proven success on these materials,<sup>35,48,49</sup> the exchange-correlation functional of Perdew-Burke-Ernzerhof (PBE)<sup>50</sup> was used in all cases and the Grimme D2 dispersion correction.<sup>51</sup> The structure was considered fully relaxed when convergence of residual forces  $< 1 \times 10^{-4} \text{ eV.\AA}^{-1}$ , and convergence of the SCF cycles  $< 1 \times 10^{-13} \text{ eV.atom}^{-1}$  was achieved. All phonons were calculated at the Brillouin zone center ( $\mathbf{k} = 0$ ) according to linear response theory, as implemented in the CASTEP suite.<sup>52</sup> The acoustic sum rule was imposed analytically. LO-TO splitting was not considered.

## Test systems

In our previous work, an up-pumping model based on sampling of the complete phonon dispersion curves was presented.<sup>35</sup> For comparison, herein we select a subset of the same systems, namely HNB (hexanitrobenzene),  $\beta$ -HMX (1,3,5,7-tetranitro-1,3,5,7-tetrazoctane),  $\alpha$ -FOX-7 (1,1-dinitro-2,2-diaminoethene),  $\alpha$ -NTO (nitrotriazolone) and TATB (1,3,5-triamino-2,4,6-trinitrobenzene). Additionally, our test set includes two new systems which crystallize in large unit cells,  $\epsilon$ -CL-20 (Hexanitrohexaazaisowurtzitane; REF: PUBMUU12) and the orthorhombic polymorph of TNT (2,4,6-trinitrotoluene; REF: ZZZMUC01), o-TNT.

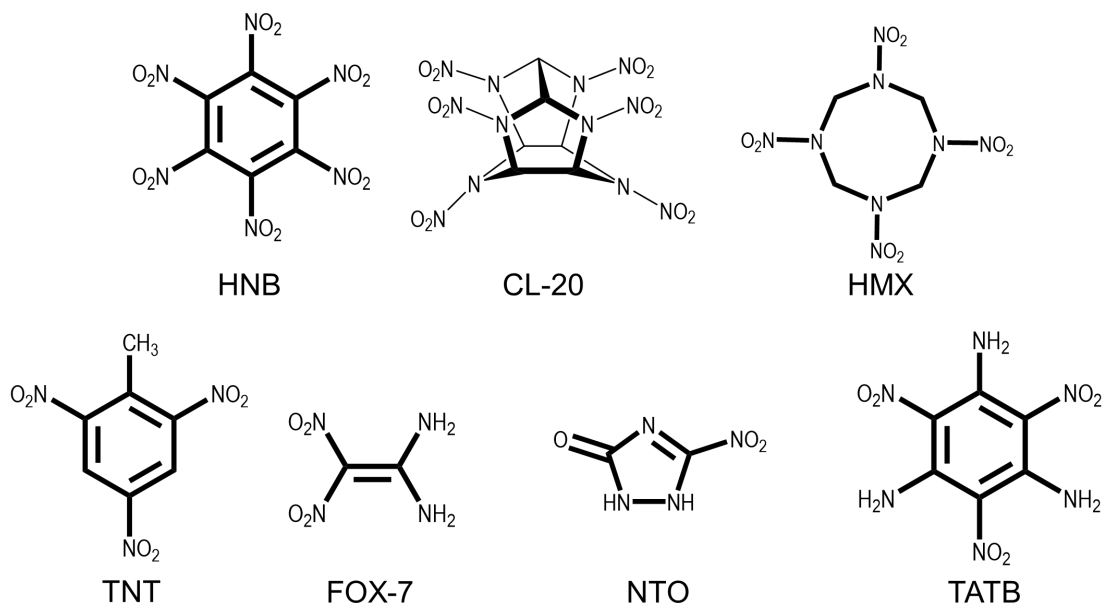


Figure 1: Schematic representation of molecular structures of the EMs used in this study.

The selection of molecules is chosen to cover a range of impact sensitivities, from the highly sensitive HNB through to the insensitive TATB. Moreover, the selected molecules represent a range of so-called ‘trigger linkages’ (*i.e.* active explosives bonds), thereby assessing the generality of the approach. The sensitivity values for each compound were taken from literature, Table 1. Our model is based on consideration of the crystallographic structure, with key parameters of our test system provided in Table 2. We note that reports of impact sensitivity values vary widely in the literature, owing to differences in experimental conditions, including particle size<sup>53</sup> and testing conditions. Where possible, a range of sensitivity values have been provided, although we note the relative ordering of compounds does not depend significantly on this variability. We note that the sensitivity of o-TNT has not been widely reported. However, with near identical crystallographic and molecular structures, density, and compressibility as compared with m-TNT,<sup>54</sup> we take its IS value to be equivalent.

Table 1: Sensitivity and selected crystallographic parameters for the EM test set. References are provided for the reported impact sensitivity values.

EM	IS /Nm	$V_{\text{exp}}/\text{\AA}^3$	$V_{\text{calc}}/\text{\AA}^3$	$\Delta V / \%$	Space Group	Ref.
HNB	2.75	581.49	585.25	+0.6	$I2/c$	55
$\epsilon$ -CL-20	3-4.25	1397.25	1461.01	+4.6	$P2_1/c$	56
$\beta$ -HMX	6.5-8	501.37	516.01	+2.9	$P2_1/c$	56,57
o-TNT	24	1823.55	1839.64	+0.8	$Pca2_1$	58
$\alpha$ -FOX-7	25.2-31.8	522.33	514.32	-1.5	$P2_1/n$	58,59
$\alpha$ -NTO	72.75	902.06	927.22	+2.8	$P\bar{1}$	57
TATB	122.5	425.25	434.69	+2.2	$P\bar{1}$	57

## Results and Discussion

### *Vibrational spectra of energetic materials*

Prior to considering a vibrational up-pumping model it is necessary to identify the upper limit of the phonon bath,  $\Omega_{max}$ . This value is intended to bound all lattice modes and must refer to a point where  $g(\omega) = 0$  in the phonon density of states (see Figure 2). While direct visualization of the eigenvectors (EV), in principle, offers a straightforward approach to marking the distinction between lattice mode and molecular mode behavior, in practice ambiguities often prevail.<sup>35</sup> In molecular crystals comprising  $Z$ ,  $N$ -atom molecules, there are six degrees of freedom (translations, rotations) for each molecule in the unit cell. It follows that each wave vector in a molecular crystal comprises  $6Z$  external modes of vibration. Within the rigid body approximation, the  $6Z$  external modes are assumed to be well separated from the remaining  $3NZ-6Z$  internal molecular modes. However, a small number of internal modes - typically rocking and wagging motions of *e.g.* NO<sub>2</sub> moieties - can often appear at frequencies that are on par with the external modes. These are known as amalgamated modes. The total number of vibrational bands which are bound below  $\Omega_{max}$  is therefore  $Z(6+Y)$ , where  $Y$  is the total number of amalgamated modes per molecule.

Table 2: Estimation of  $\Omega_{max}$  based on the mode-counting (MC), eigenvector mode-counting (EV-MC), centre of mass (CoM), and heat capacity (HC) approaches. The number of molecules in the unit cell,  $Z$ , number of atoms per molecule,  $N$ , and the number of amalgamated modes per molecule,  $Y$ , are given.

EM	$Z$ / unit cell	$N$ /molecule	$Y$	$\Omega_{max}/\text{cm}^{-1}$			
				MC	EV-MC	CoM	HC*
HNB	2	24	9	215	140	140	146
$\epsilon$ -CL-20	4	36	16	217	217	217	178 / 222
$\beta$ -HMX	2	28	11	243	191	191	100 / 198
o-TNT	8	21	9	212	212	212	160 / 217
$\alpha$ -FOX-7	4	14	3	169	169	169	174
$\alpha$ -NTO	8	11	2	254	199	199	166 / 204
TATB	2	24	6	140	140	140	144

\* Multiple values reported due to onset of plateau regions in HC model below the MC- $\Omega_{max}$  value.

The first approach we adopted to identify an appropriate value for  $\Omega_{max}$  was mode-counting (MC), wherein a direct comparison between the isolated molecule vibrational spectrum and its corresponding condensed phase density of states was made. The lowest frequency  $6Z$  external modes can be readily identified in the condensed phase calculation, thereby defining a maximum cut-off frequency for the external mode frequencies,  $\Omega_{ext}$ . From our gas phase simulations, the number of internal modes which fall below  $\Omega_{ext}$  are subsequently identified,  $N_{int}$ . Owing to the symmetry splitting of molecular modes in the solid state, this yields  $Y = ZN_{int}$  amalgamated modes. As these  $Y$  amalgamated modes do not contribute to the required  $6Z$  external modes, we extend the limit of  $\Omega_{ext}$  to include the additional  $Y$  vibrational frequencies. Additional molecular modes which fall below the extended  $\Omega_{ext}$  are included, and the value of  $\Omega_{ext}$  is again updated. This procedure is repeated until no further molecular modes can be identified below the  $6N$  external mode maximum limit. The final cut-off is taken to be the top of the phonon bath,  $\Omega_{max}$ . Values of  $\Omega_{max}$  identified using this MC procedure are tabulated in Table 2 and discussed further in the ESI. It is worth highlighting that the values of  $\Omega_{max}$  obtained in this way are generally consistent with our previous report based on direct visualization of the EVs.<sup>35</sup> In a few cases (notably, HNB,  $\beta$ -HMX and  $\alpha$ -FOX-7), crystal field effects lead to marked shifts in the gas phase frequencies. Where this is the case, EV visualization can then be invoked to facilitate the tuning of MC-derived  $\Omega_{max}$  values (see Table 2 and ESI). In practice, the absolute



eigenvalues obtained for both MC- and EV-MC- $\Omega_{max}$  are extended by  $5 \text{ cm}^{-1}$  when applied to the Gaussian-broadened phonon density of states.

To facilitate EV analysis, we note that internal modes should be characterized by negligible displacement of the molecular centre of mass (CoM), which contrasts with the marked displacement of CoM expected for external modes. Correspondingly, we explored the possibility for rapid EV analysis – and hence identification of  $\Omega_{max}$  – by tracking the variation in CoM as expressed by the  $3N$  rectilinear eigenvectors. In this way we obtain  $\Omega_{max}$  values as the absolute vibrational eigenvalues above which no eigenvectors display CoM displacements (see ESI). In practice  $5 \text{ cm}^{-1}$  are again added to these values to allow for Gaussian broadening. The results obtained are essentially identical to the EV-MC method, which was as expected as both involve direct analysis of the eigenvectors. This route does however provide a quantitative analysis that is much less subjective than the visual comparison inherent in the EV-MC method and also provides a route to the automated screening for  $\Omega_{max}$  values.

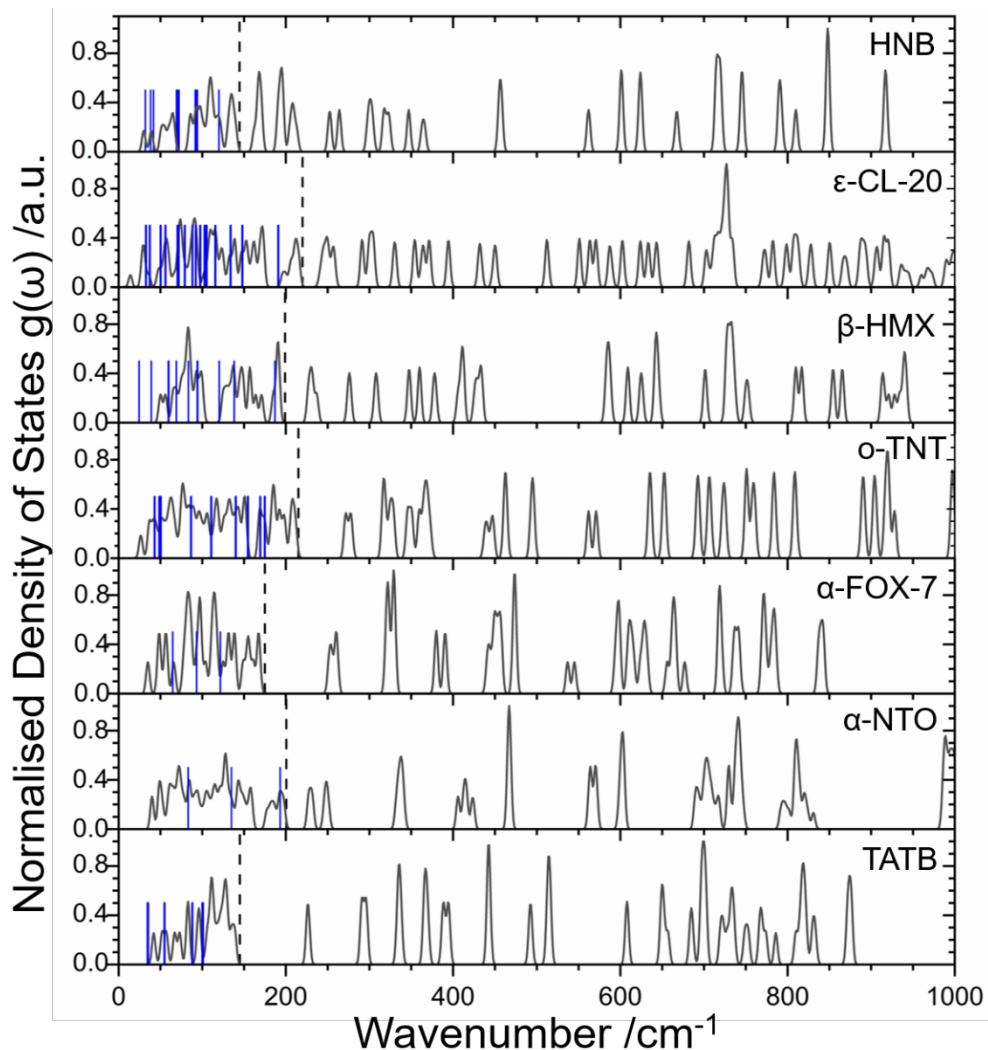


Figure 2: Phonon density of states for EMs in the test set.  $\Omega_{max}$  as identified by the EV-MC and CoM methods is indicated with a dashed vertical line, and the gas phase frequencies below this point are shown as blue vertical lines.

With the limiting value  $\Omega_{max}$  set by the EV-MC (or analogously CoM) method, it becomes straightforward to calculate the heat capacity values required to solve Equation 3. Based on the Brillouin zone center vibrational frequencies only, an approximate density of states was generated by applying a Gaussian smearing of  $\pm 5 \text{ cm}^{-1}$ . This yields a FWHM for isolated vibrational bands of  $\pm 5 \text{ cm}^{-1}$ .<sup>42</sup> We therefore obtain the bulk heat capacity by Equation 3, and the phonon bath heat capacity *via* the same equation by imposing an upper integration limit of  $\Omega_{max}$ , see ESI. It follows from Equation 3 that the adiabatic heating in each case becomes proportional to the ratio  $C_{tot}/C_{ph}$ . The high temperature limit of the bulk and phonon heat capacities can then be determined according to,

$$c_v = \int_0^{\Omega_{max}} \frac{\partial n}{\partial T} \hbar \omega g(\omega) d\omega \quad \text{Eqn 7}$$

where  $\omega_i = \Omega_{max}$  for the phonon heat capacity and  $\omega_i$  extends across all  $3N$  vibrations for the total heat capacity. Based solely on the zone center vibrational frequencies, it is worth noting that the heat capacities are expected to be over-representative for each system (see ESI). However, as all subsequent properties will be calculated from the same zone-center frequencies, we expect this to be intrinsically consistent with the remainder of the model.

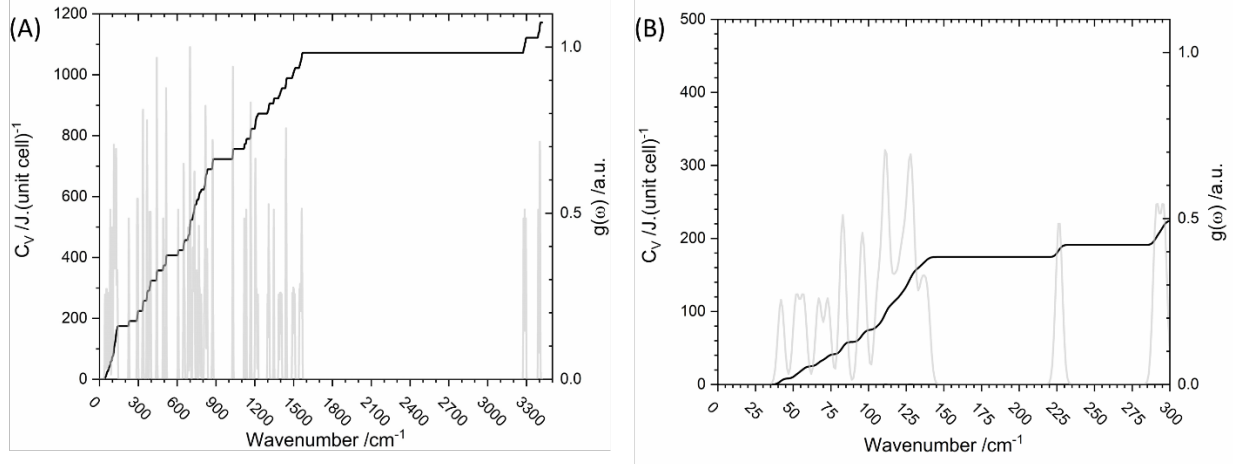


Figure 3: Calculation of vibrational heat capacity for TATB based on phonon density of states, according to Eqn 7. (A) Overlay of cumulative vibrational heat capacity and phonon density of states. (B) Identification of  $\Omega_{max}$  based on plateau in cumulative heat capacity.

It is worth highlighting that visualization of the cumulative total vibrational heat capacity, Figure 3, with respect to increasing wavenumber provides another convenient way to visually locate  $\Omega_{max}$ . Our assumption for  $\Omega_{max}$  requires  $g(\omega) = 0$ , which is given by a plateau in the cumulative heat capacity curves, Figure 3b. If the first such plateau is taken to be indicative of  $\Omega_{max}$  – denoting the first break in  $g(\omega)$  – reasonable agreement with our MC approach is obtained, Table 2. In many cases, however, a small plateau is observed at wavenumbers well below predicted values of  $\Omega_{max}$  by the MC method. To understand the sensitivity of our up-pumping model to the choice of  $\Omega_{max}$  we will consider the potential of these lower bounds; note, however, the lower values generally do not bound the number of modes which are amalgamated by the MC, EV-MC or CoM approaches (see ESI, Table 2).

### Predicting impact sensitivity

A critical aspect of the up-pumping model, described above, pertains to normalization of the phonon density of states,  $g(\omega)$ . Although a count of  $g(\omega)$  is provided by the DFT simulations, it accounts only for pathways which reside at  $\mathbf{k}=0$ , or from an arbitrary number of  $\mathbf{k}$  points. To reflect better the true relative number of vibrational pathways present in each system, we apply a normalization of  $\int^{\Omega_{\max}} d\omega g(\omega) = \frac{Z(6+Y)}{V_0}$ . A solid with volume  $V$  and containing  $N$  unit cells comprises  $N$  wave vectors. Hence in a region  $\mathbf{k}+d\mathbf{k}$  the number of wave vectors (described by Equation 8) differs between systems only by the volume of the unit cell.

$$g(\mathbf{k})d\mathbf{k} = \frac{V}{2\pi^3} 4\pi k^2 d\mathbf{k} \quad \text{Eqn 8}$$

It therefore follows that normalizing by the unit cell volume offers a direct route for comparing the relative number of transfer pathways available across different EM crystalline solids.

To reflect the relative amount of energy transferred to each EM upon impact we assume that the compressibility factor and Grüneisen parameters expressed in Equation 1 are constant across all EMs (see ESI). This is generally a good approximation based on experimentally available data.<sup>60</sup> For the purpose of this work, we take as a model a value of  $V/V_0 \approx 0.8$ , with a value for  $\gamma \approx 4$ . This results in a final temperature,  $T_f = 953.6$  K, and hence a total bulk heating  $\Delta T = 655.6$  K. For each material, the corresponding initial quasi-temperature of the phonon path,  $\phi(0)_{ph}$ , is obtained by Equation 7, and is hence directly proportional to  $C_{tot}/C_{ph}$ , Table 3 (and ESI).

Table 3: Predicted total and phonon-bath only heat capacities, and phonon bath shock temperature heating used in the IS prediction model, calculated based on EV-MC and HC derived values of  $\Omega_{\max}$  (shown in Table 2; lowest values employed in the HC-model to explore the sensitivity of up-pumping to  $\Omega_{\max}$ ).

	EV-MC/CoM				HC Model		
	$C_{tot}$	$C_{ph}$	$C_{tot}/C_{ph}$	$\phi(0)_{ph}/K$	$C_{ph}$	$C_{tot}/C_{ph}$	$\phi(0)_{ph}/K$
HNB	585.85	112.20	5.22	3423.3	112.20	5.22	3423.3
$\epsilon$ -CL-20	892.89	178.69	5.00	3278.0	157.92	5.65	3704.1
$\beta$ -HMX	685.48	128.82	5.32	3487.8	62.33	11.00	7211.6
o-TNT	521.14	122.59	4.25	2786.3	89.35	5.83	3822.1
$\alpha$ -FOX-7	342.76	68.56	5.00	3278.0	68.56	5.00	3278.0
$\alpha$ -NTO	255.45	63.33	4.03	2642.1	55.03	4.64	3042.0
TATB	585.78	87.26	6.71	4399.1	87.26	6.71	4399.1

Following normalization of the phonon density of states, our up-pumping model follows the two-step process described in Equations 6. At each step of the up-pumping calculation, the up-pumped states are dispersed over the  $Z$  molecules, reflecting the localization of the vibrational energy with increasing frequency. A shock temperature (Table 3) is initially applied to the underlying phonon DOS (pre-set with populations corresponding to 300 K), yielding a visibly ‘hot’ phonon bath, Figure 4a. Following from Equation 6a, this excess energy is propagated onto the states residing at  $\Omega_{\max} < \omega < 2\Omega_{\max}$ , Figure 4b. Finally, Equation 6b allows for the final up-pumping of density to a maximum of  $\omega < 3\Omega_{\max}$ , Figure 4c.

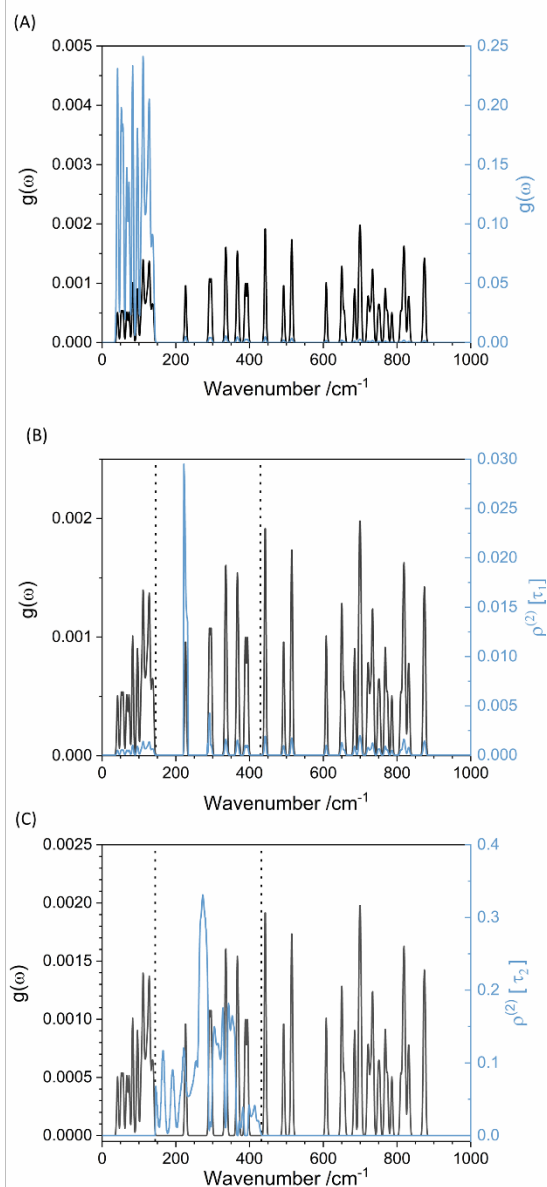


Figure 4: Key steps in the two-step up-pumping model for TATB. (A) Hot phonon bath and thermally equilibrated internal modes; (B) Hot doorway states obtained by up-pumping through Equation 6a; (C) Up-pumped density obtained *via* Equation 6b and projected onto fundamental bands. In each case, the phonon density of states  $g(\omega)$  is shown in black, with the up-pumping effects on the populations shown in blue. In (B) and (C) the vertical dotted lines denote  $\Omega_{\max}$  and  $3\Omega_{\max}$ .

Propagation of vibrational density in this way provides a measure of the total excitation of the system. However, to yield initiation, the vibrational energy must induce an electronic response. This process is known as dynamical metallization, wherein heightened vibrational excitation of the target frequency induces structural changes that result in a closing of the electronic band gap.<sup>37</sup> In simple systems, such as metal azides, there exists a well-defined ‘target’ vibrational mode that clearly facilitates the necessary metallization phenomenon. When considering more complex molecules, as done here, no single vibrational mode can be readily identified as being responsible for initiation. Instead, we adopt the arguments of RRKM

kinetic theory of unimolecular decomposition. Noting that for large molecules the two-phonon density of states is approximately continuous over the range  $\Omega_{max} - 3\Omega_{max}$  (see ESI), all up-pumped energy into this region can rapidly redistribute into the underlying fundamental modes present in the original density of states. Under this assumption, the probability of unimolecular decomposition is proportion to the total amount of up-pumped vibrational energy. Correspondingly, we take the integrated density of activated fundamental states as being indicative of the impact sensitivity, Figure 5.

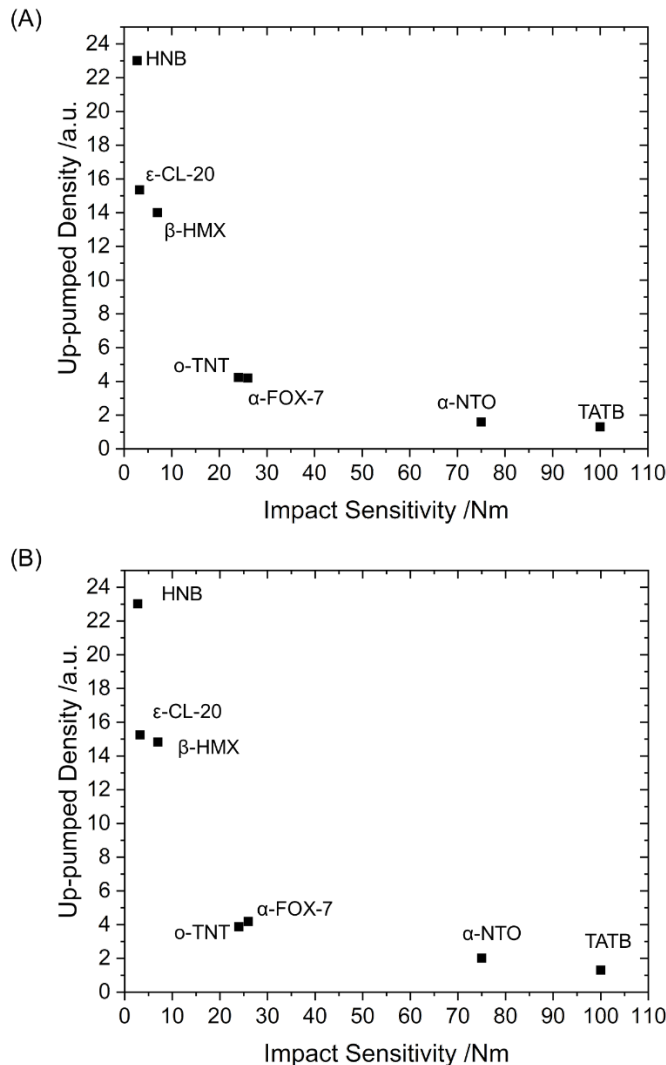


Figure 5: Predicted relative sensitivity using the vibrational up-pumping model. (A) Based on EV-MC/CoM- $\Omega_{max}$  values; and (B) Based on lowest HC- $\Omega_{max}$  values. Further details are available in the ESI, along with predicted sensitivities based on the MC- $\Omega_{max}$  values.

The results in Figure 5A clearly indicate that our up-pumping model using the combined EV-MC or CoM- $\Omega_{max}$  values captures the relative experimental IS ordering for our EM test set well. Whereas negligible vibrational population is up-converted in the insensitive compounds  $\alpha$ -NTO and TATB, the sensitive

compounds HNB and  $\epsilon$ -CL-20 exhibit large up-pumping contributions. Moreover, the ordering of the intermediate compounds  $\beta$ -HMX, o-TNT, and  $\alpha$ -FOX-7 is also correctly predicted. The correct ordering of sensitive compounds HNB and  $\epsilon$ -CL-20 – which differ by only 0.25 Nm – is correctly captured, although a minor mis-ordering of o-TNT and  $\alpha$ -FOX-7 ( $\Delta$ IS *ca* 1 Nm) is observed. We note, however, that such small differences in IS values are not generally meaningful, and are well within the errors of the experimental testing approach. Even the use of the HC- $\Omega_{max}$  values provides a reasonable ordering outcome (Figure 5B), with the most notable change being a reduced resolution in the relative predicted sensitivities for  $\beta$ -HMX and  $\epsilon$ -CL-20. Despite these minor losses in the predictive capacity of the model, the overall classification of the test materials is, however, retained.

Our data therefore suggest that, while the choice of  $\Omega_{max}$  is undoubtedly important in predicting IS, it should be thought of as an indirect parameter only. From a materials design perspective, this also suggests that the initial material excitation (*i.e.* the value of  $\phi_{ph}$ ) is not directly indicative of the impact reactivity of the EMs tested here. This is further evident from the tabulated values of  $\phi_{ph}$  in Table 3, and a lack of direct correlation with impact sensitivity response. Correspondingly, features which directly influence the bulk and phonon heat capacities (*c.f.* Equation 3), are unlikely to provide sensitive design targets for new EMs. For example, a low phonon shock quasi-temperature is expected for systems comprising simultaneously a low total heat capacity (induced by a small number of molecular degrees of freedom), and a high phonon heat capacity (induced by a high  $Z$  count in the primitive unit cell and a large number of  $Y$  amalgamated molecular vibrational modes). While these variables are undoubtedly important in understanding the vibrational up-pumping process, we must look to other areas to influence the structure/property relationship.

The predictions obtained from our model are instead dominated by the underlying vibrational structure. In particular, the model is dominated by the availability of vibrational states in the range  $\Omega_{max} < \omega < 2\Omega_{max}$ ,<sup>35</sup> *i.e.* the region which captures directly the energy up-pumped from the ‘hot’ phonon bath, Equation 6a. This dominant aspect of the up-pumping model has been noted previously in the literature.<sup>38</sup> The up-pumped density in the region  $2\Omega_{max} < \omega < 3\Omega_{max}$  captures the second stage of the up-converted phonon energy, Equation 6b, and also plays a crucial role in determining the predicted sensitivity ordering. It is therefore likely that influencing the distribution of vibrational states within these windows will provide the design-space for new energetic materials with tailored impact sensitivity responses.

The range of vibrational modes driving the phonon up-pumping model spans *ca.*  $200 \pm 50 \text{ cm}^{-1}$  to  $600 \pm 150 \text{ cm}^{-1}$ . This is the vibrational window determined largely by the flexibility of the energetic molecule, including angle bends, ring deformations, and other similar molecular distortions. Correspondingly, it is reasonable to expect that molecular flexibility will be indicative of the vibrational density in this critical window, and hence of the overall sensitivity of the material. To this end, we calculated two established flexibility indices: (1) the rotatable bond count (*i.e.* the number of single bonds which are not in a ring and which are bound to a non-terminal heavy atom), and (2) the Kier Molecular Flexibility (KMF) index, given by<sup>61</sup>

$$\text{KMF} = \frac{K_{\alpha}^1 K_{\alpha}^2}{N} \quad \text{Eqn 9}$$

where  $K_{\alpha}^1$  and  $K_{\alpha}^2$  represent the first- and second-order shape indices (*i.e.* molecular graph parameters indicative of the atom (vertex) count, vertex cycles, and the degree of branching) and  $N$  is the number of vertex points in the molecule. Both molecular flexibility indices yield very similar data to each other and, with a few anomalous data points, provide a generally good indication of the relative sensitivity of the EMs in our test set, Figure 6. The flexibility indices of TATB are on par with TNT, despite a pronounced

difference in their IS response. Similarly, CL-20 exhibits a KMF that is nearly two orders of magnitude above the remaining EMs, although overall it is correctly predicted to be highly sensitive. It follows that designing energetic molecules with a high degree of molecular flexibility is likely to be a strong indication of overall sensitivity. This is consistent with our up-pumping model, where this flexibility is encoded into the molecular vibrations that fall into the  $\Omega_{max} < \omega < 3\Omega_{max}$  window. This same flexibility can be influenced by the strength of interatomic interactions, which can hinder rotations and modify charge densities and bond dissociation energies. Thus, tuning intermolecular interactions, through polymorphism or multi-component crystallization can also be expected to tune reactivity to mechanical impact.

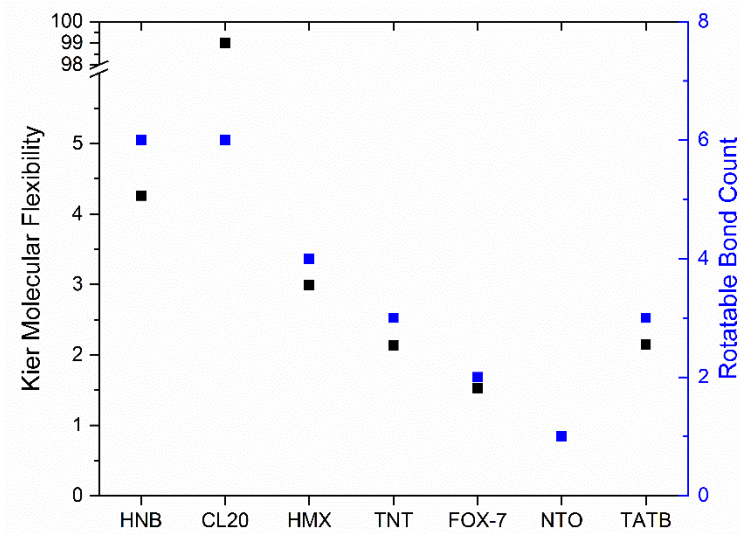


Figure 6: Comparison of molecular flexibility parameters against impact sensitivity ordering, listed (left to right) in order of decreasing impact sensitivity.

### Summary of up-pumping approach

In summary, the proposed up-pumping model for IS prediction follows five principle steps, which are expected to be easily automated: (1) Following simulation of zone centre vibrational frequencies, we elect to impose a Gaussian broadening of  $5 \text{ cm}^{-1}$  to approximate phonon dispersion. The Gaussian broadened spectrum is normalized by the total number of wave vectors in the material *via* Equation 8. This provides the underlying vibrational spectrum upon which up-pumping occurs. (2) The initial conditions for up-pumping require the definition of a phonon bath, with upper bound  $\Omega_{max}$ . Our results suggest that  $\Omega_{max}$  can be readily obtained by analyzing the displacement of the molecular centre of mass (CoM) for each simulated eigenvector. This offers a clear indication of the frequency which bounds all external phonon modes. (3) The quasi-temperature of the phonon bath is subsequently obtained by considering adiabatic compression, Equation 1, in which all EMs are here assumed to have identical Grüneisen coefficients and compressibility. Correspondingly, the final bulk temperature of each material is equivalent. By calculating the total and phonon heat capacities, Equation 7, the material-specific phonon quasi-temperature is obtained *via*  $C_{tot}/C_{ph}$ . (4) This excess phonon population is up-pumped *via* Equations 6a and 6b. (5) Finally, as our approach is limited to first order anharmonic effects, we take the total quantity of up-pumped population within the region  $\Omega_{max} < \omega < 3\Omega_{max}$  as being indicative of the relative IS of the material.

### Conclusions

The search for new EMs that offer the correct balance of power, safety, and compliance with environmental legislation is an inherently hazardous endeavor. Herein we have demonstrated that impact sensitivity, one of the key safety metrics for EMs, is amenable to prediction via a two-stage vibrational up-pumping model based on the Brillouin zone-centre vibrational density of states. Our model correctly predicts the relative impact sensitivities for a test set of molecular crystals that are both structurally and energetically diverse and are some of the most well-known EMs in commercial and academic research fields. Moreover, this is achieved without resort to any fitting parameters, meaning that it can be applied to any energetic material, provided the crystal structure is known. As a critical decision in implementing the up-pumping approach, we have explored a series of possible selections for  $\Omega_{max}$ , namely a mode counting approach (MC), its combination with manual eigenvector analysis (EV-MC), a related approach based on screening displacements of molecular centre of mass (CoM), and analysis of cumulative heat capacities (HC). Although we expect the EV-MC approach to be most physically meaningful, it requires significant manual efforts, and is therefore not amenable to automated screening. The CoM approach provides the same values of  $\Omega_{max}$  as EV-MC and is readily implemented for automated and high-throughput screening. We therefore expect the CoM methodology to become useful for screening libraries of potential EMs.

The up-pumping model goes beyond simply reproducing known experimental data. Having a predictive tool allows exploration of structure/property relationships at the most fundamental of levels, and thereby provides a feedback loop to rationalize material design with tailored energetic properties. Here the distribution of molecular modes that reside within the  $1-3\Omega_{max}$  range hold the key to designing a material that is classed as either a primary (highly impact sensitive) or a secondary (lower impact sensitive) energetic. It therefore follows that molecular design, through the introduction (or removal) of low-energy amalgamated vibrational modes, and studies on polymorphism and multi-component crystallization, to increase (or reduce) the strengths of intermolecular interactions and charge distributions through the lattice all have a part to play in shifting the pattern and distribution of low energy molecular vibrations. This, in turn, affects the efficiency of the molecular crystal to adsorb the energy from the mechanical impact. The vibrational up-pumping model draws a direct line from the IS prediction models based on electronic structure through to a physical model of energy transfer and bond dissociation. It is clear that having reliable computational modelling tools, to act in a screening capacity in the search for new energetic materials, is an important step forward in the field of energetic material research.

### **Data availability**

The data that supports the findings of this study are available within the article (and its supplementary material).

### **Supplementary material**

See supplementary material for the simulated crystallographic structures, computed vibrational spectra, and further details on the calculation of solid-state heat capacities and vibrational up-pumping processes.

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